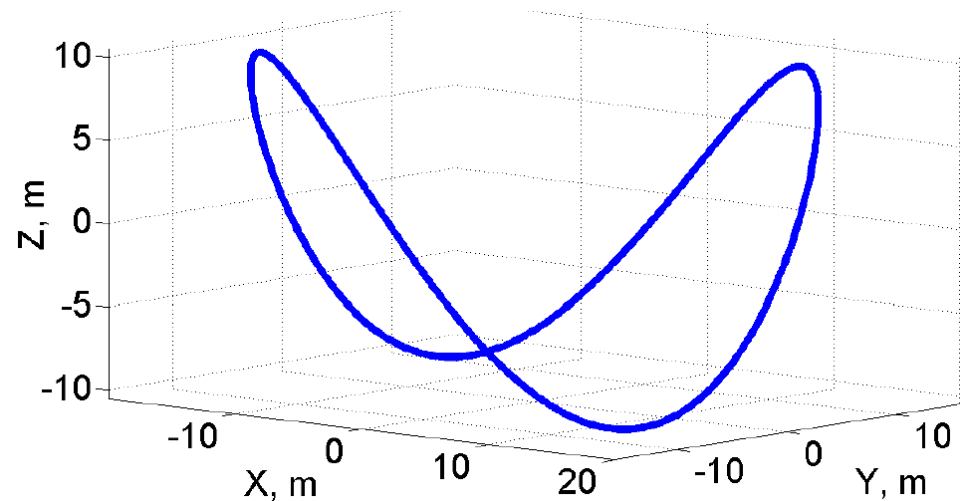
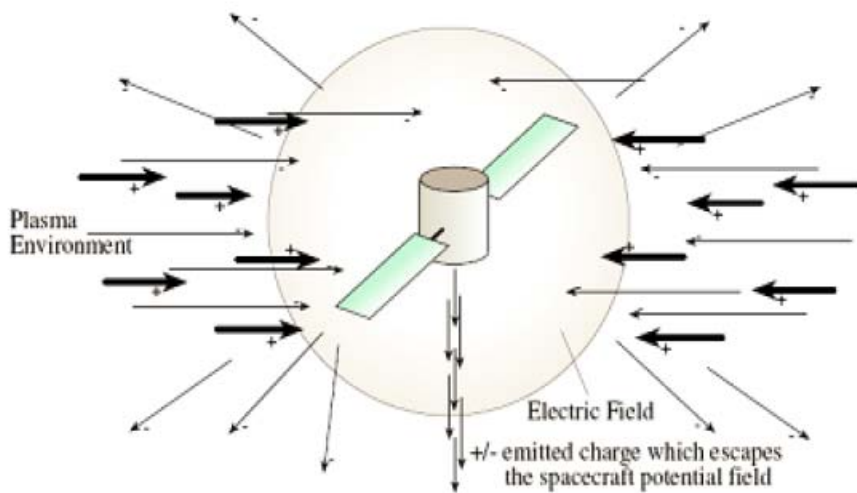


# Periodic Relative Orbits of Two Spacecraft Subject to Differential Gravity and Coulomb Forces

Presented by Daan Stevenson for  
Drew R. Jones (author) & Hanspeter Schaub (co-author)

5<sup>th</sup> SFFMT Conference

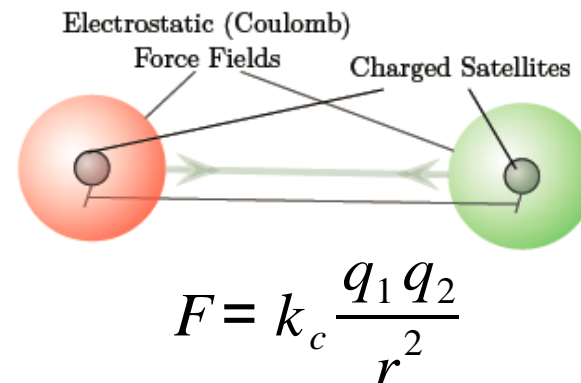
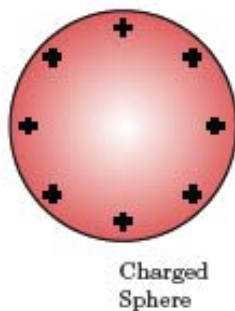
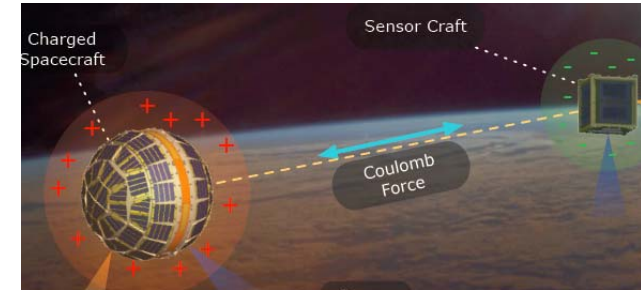
May 31, 2013



# Coulomb Formation Motivation

## Coulomb Formation: Close-flying charged craft using electrostatic forces

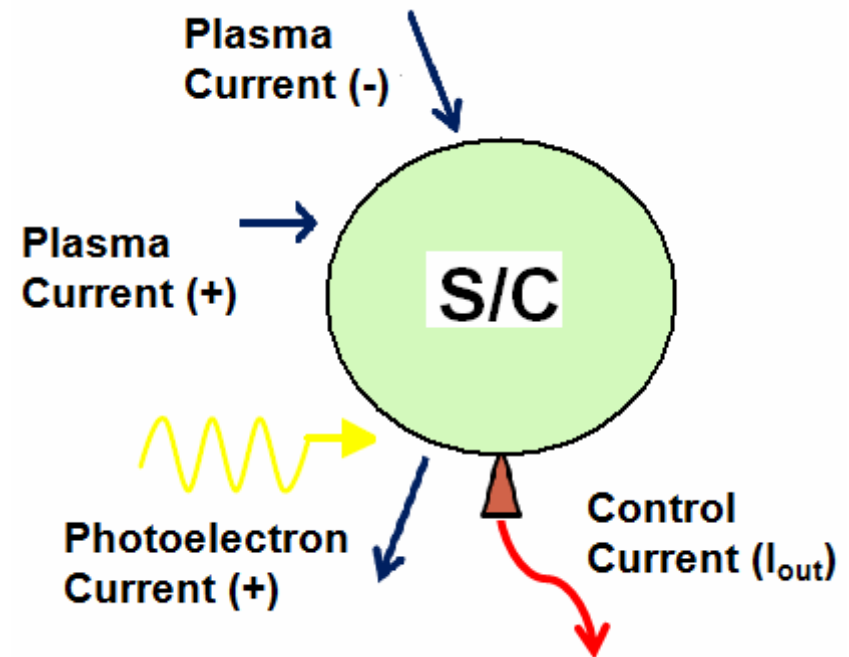
- Very efficient (ISP~ $10^{13}$  s)
- Forces have limited reachability
- Forces introduce NL coupling
- Forces complex to model in the plasma



- I. Approximate s/c charging and Coulomb force model
- II. Dynamics and dynamical system properties
- III. Derivation of periodic Coulomb formation solutions
- IV. Simulated periodic Coulomb formation solutions
- V. Perturbed Coulomb formation motions
- VI. Conclusions and future research

# Spacecraft Charging

- Spacecraft naturally assume non-zero potential (charge)
- Artificially altering potential has been demonstrated
- Spherical s/c model adopted
  - Somewhat an abstraction
  - Enables 1<sup>st</sup> order calculations and assumptions

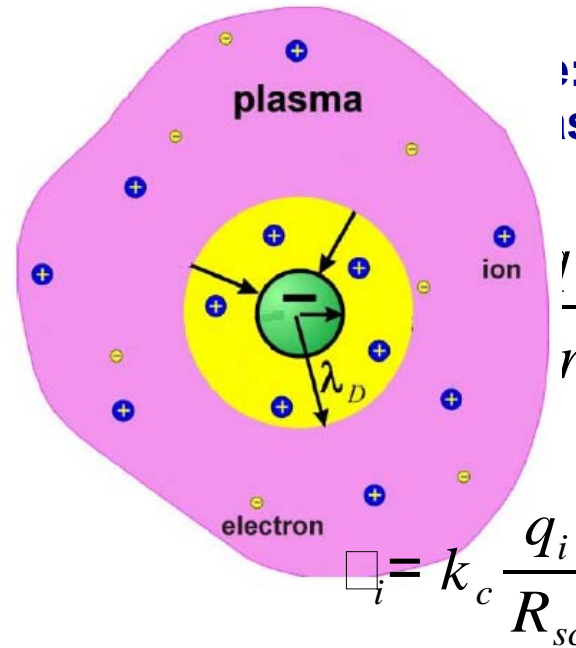


# Electrostatic Model

- Spacecraft are not point charges in a vacuum...
  - Adopt approximate and analytical Coulomb force model
  - But still account for plasma shielding effect
  - Assume formations near GEO

**Coulomb Force: Point Charges in Vacuum**

$$F_{12} = k_c \frac{q_1 q_2}{r_{12}^2}$$

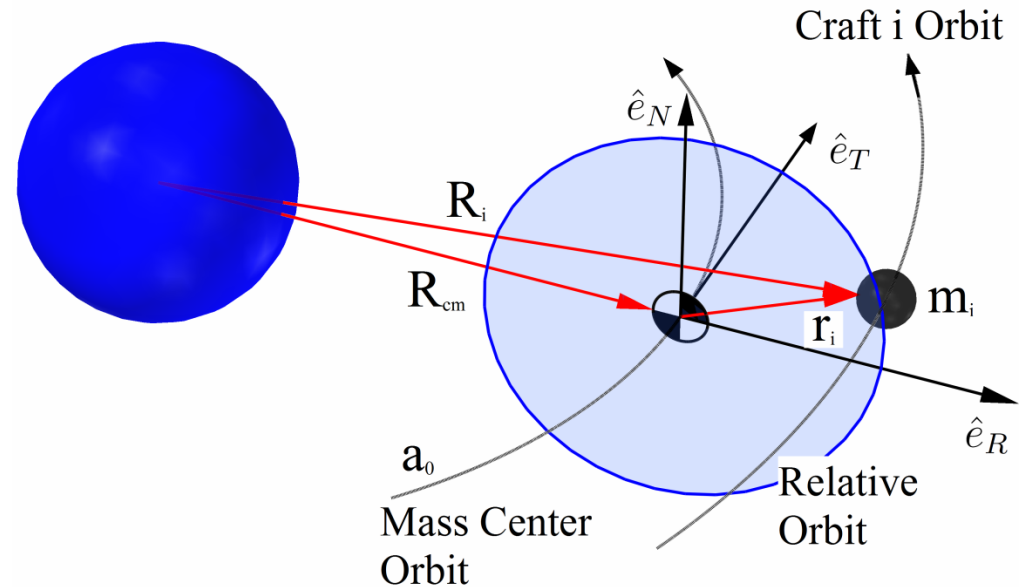


**Finite Bodies with Plasma Shielding**

$$F_{12} = \frac{q_1 q_2 (1 + r_{12}/\lambda_d)}{r_{12}^2 \exp[r_{12}/\lambda_d]}$$

- Hill Frame Model
- Clohessy-Wiltshire relative motion EOM

$$\sum_i m_i \mathbf{r}_i = 0$$



- With Debye-Hückel Coulomb force model:

$$\ddot{\mathbf{r}}_i = \begin{bmatrix} 2\omega\dot{y}_i + 3\omega^2 x_i \\ -2\omega\dot{x}_i \\ -\omega^2 z_i \end{bmatrix} + \left[ \frac{k_c q_i}{m_i} \sum_{j \neq i} \frac{q_j e^{-r_{ij}/\lambda_d}}{r_{ij}^3} \left( 1 + \frac{r_{ij}}{\lambda_d} \right) \mathbf{r}_{ij} \right]$$

- Dynamics normalized, one craft removed
  - s/c 2 motion is always dependent, via CM constraint
  - Scaled charge products introduce time transformation

$$Q_{12} = q_1 q_2$$
$$\tilde{Q}_{12} = \frac{k_c Q_{12}}{\omega^2} \quad d\tau = \omega dt \quad (\zeta)' = \frac{d\zeta}{d\tau} = \frac{1}{\omega} \frac{d\zeta}{dt}$$

- Scalar constant of motion is shown to exist
  - Holds for N-craft Hill frame systems with internal forcing
  - For two-craft formation, the motion constant is:

$$x(\tau)y''(\tau) - y(\tau)x''(\tau) = 0$$

# Dynamical System Analysis

- Seek periodic solutions to vector ODE  $\mathbf{X}' = \mathbf{F}(\mathbf{X}, \mathbf{u}, \tau)$ 
  - Dynamics linearized about such solutions

$$\delta \mathbf{X}'(\tau) = \left. \left( \frac{\partial \mathbf{F}}{\partial \mathbf{X}} \right) \right|_{(\mathbf{X}^*, \mathbf{u}^*)} \delta \mathbf{X}(\tau) = \mathbf{A}(\tau) \delta \mathbf{X}(\tau) \quad \delta \mathbf{X}(\tau) = \Phi(\tau, 0) \delta \mathbf{X}(0)$$
$$\Phi'(\tau, 0) = \mathbf{A}(\tau) \Phi(\tau, 0)$$

- STM matrix has identical symplectic property as known for linear dynamics about libration points in CRTBP
- Asymptotically stable solutions cannot exist

$$\mathbf{J}\mathbf{A}^T = -\mathbf{A}\mathbf{J} \quad \Phi\mathbf{J}\Phi^T = \mathbf{J} \quad \mathbf{J} = \left[ \begin{array}{c|c} \mathbf{0} & \mathbf{I} \\ \hline -\mathbf{I} & \mathbf{G} \end{array} \right]$$

1.  $\det(\Phi) = |\Phi| = 1$

2. At least one Floquet multiplier has modulus of unity:  $|\sigma_i| = 1$

3. The  $\sigma_i$  appear in reciprocal pairs (i.e. if  $\sigma_i$  is eigenvalue, then so is  $\sigma_j = 1/\sigma_i$ )



## Periodic Coulomb Formations

- Assume positions as simple harmonic oscillators
  - Solve time-varying charges that produce assumed motion
  - Motions are restricted (cannot be assumed arbitrarily)
- Relative instability measured via Floquet multipliers
- Dynamical coupling permits 3 solution families
  - Orbit-Normal oscillations only (1D)
  - Reference orbit-plane (2D) motions
  - Full state (3D) motions

# Reference Orbit Plane Solutions

$$x(\tau) = A_x \cos(\theta_x \tau) \quad y(\tau) = A_y \sin(\theta_y \tau)$$

$$1) \quad x(\tau)y''(\tau) - y(\tau)x''(\tau) = 0$$

$$\longrightarrow \quad \theta_x = \theta_y$$

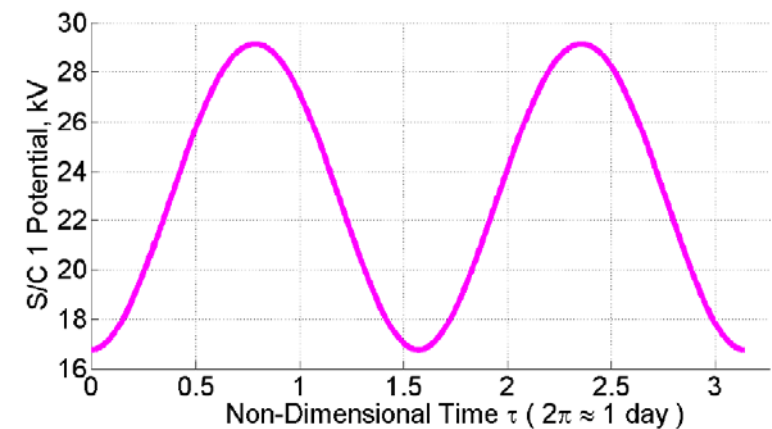
$$2) \quad \text{Dynamical Coupling} \\ \text{(Coulomb terms):}$$

$$\longrightarrow \quad \left( \frac{A_y}{A_x} \right) = \frac{-3 \pm \sqrt{9 + 16\theta^2}}{4\theta}$$

- Two roots of quadratic yield two families (cases A/B)

- Relative orbit is ellipse
- Major axis along either radial (case A) or transverse (case B)
- Short period (case B) orbits  $\rightarrow$  least unstable

$$\tilde{Q}(r(\tau)) = \frac{-1}{\Psi(r)} \left[ \theta^2 + 3 + \left( \frac{-3 \pm \sqrt{9 + 16\theta^2}}{2} \right) \right]$$



(b) Case B:  $A_x < A_y$

# Full State Solutions

- Same in-plane assumed functions (same restrictions)

- Additionally

$$z(\tau) = A_z \sin(B_z \theta \tau)$$

- 1) **Out of plane frequency factor must be even integer**



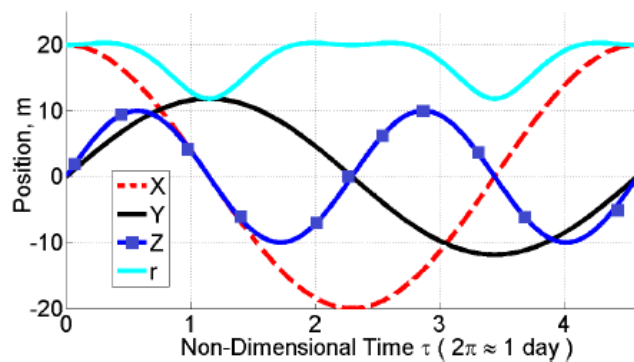
$$B_z = 2, 4, 8 \dots$$

- 2) **Time period no longer free, must satisfy:**

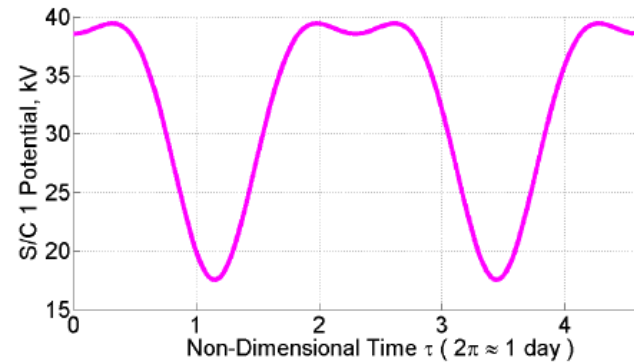


$$8\theta^2 + \left(-3 \pm \sqrt{9 + 16\theta^2}\right) [\theta^2(1 - B_z^2) + 1] = 0$$

- Projected trajectory is ellipse aligned as in case A/B
- Case B orbits have smaller max. Floquet multiplier



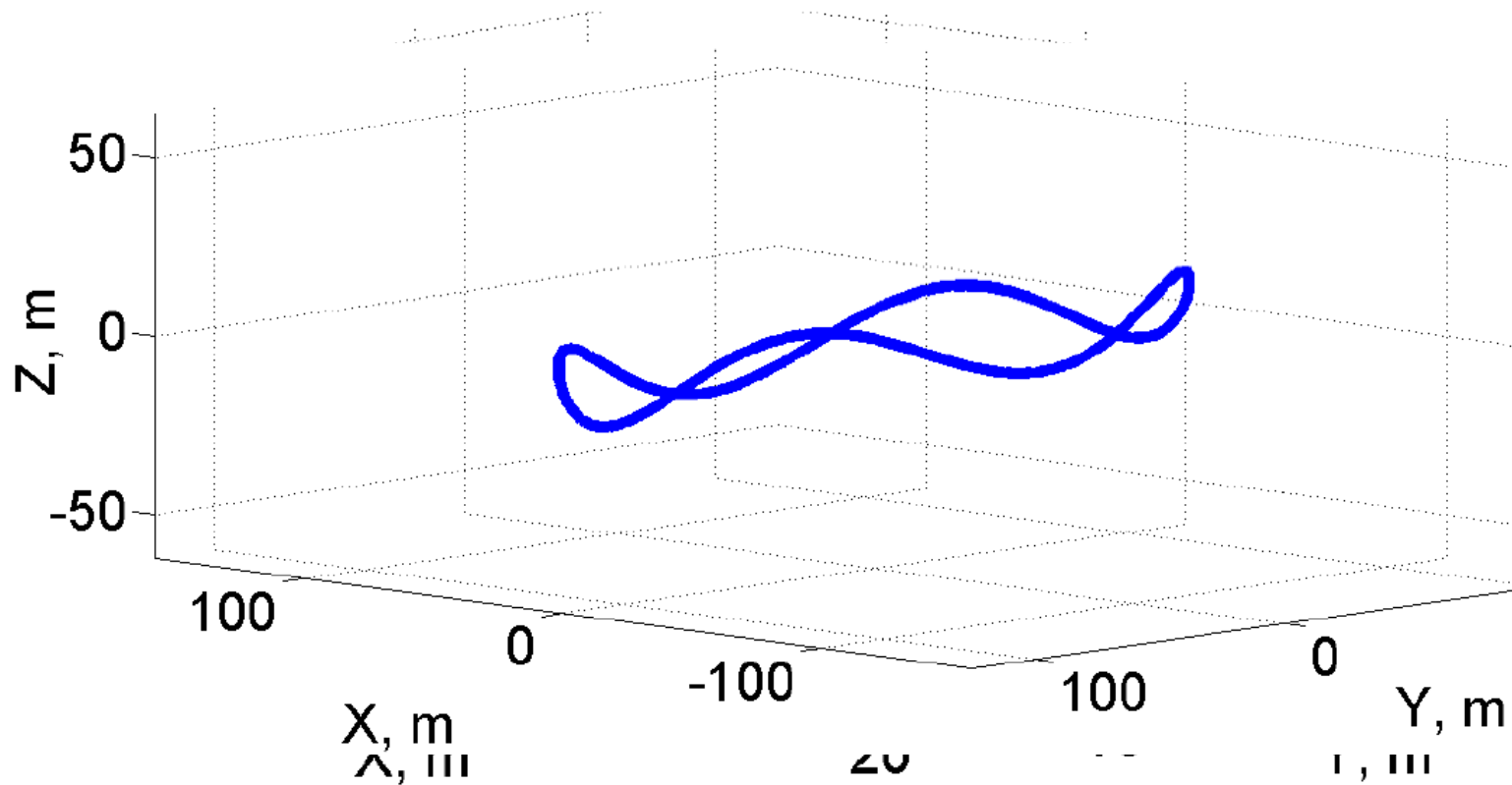
(a) S/C 1 Position History



(b) S/C 1 Potential History

# Periodic Coulomb Formations

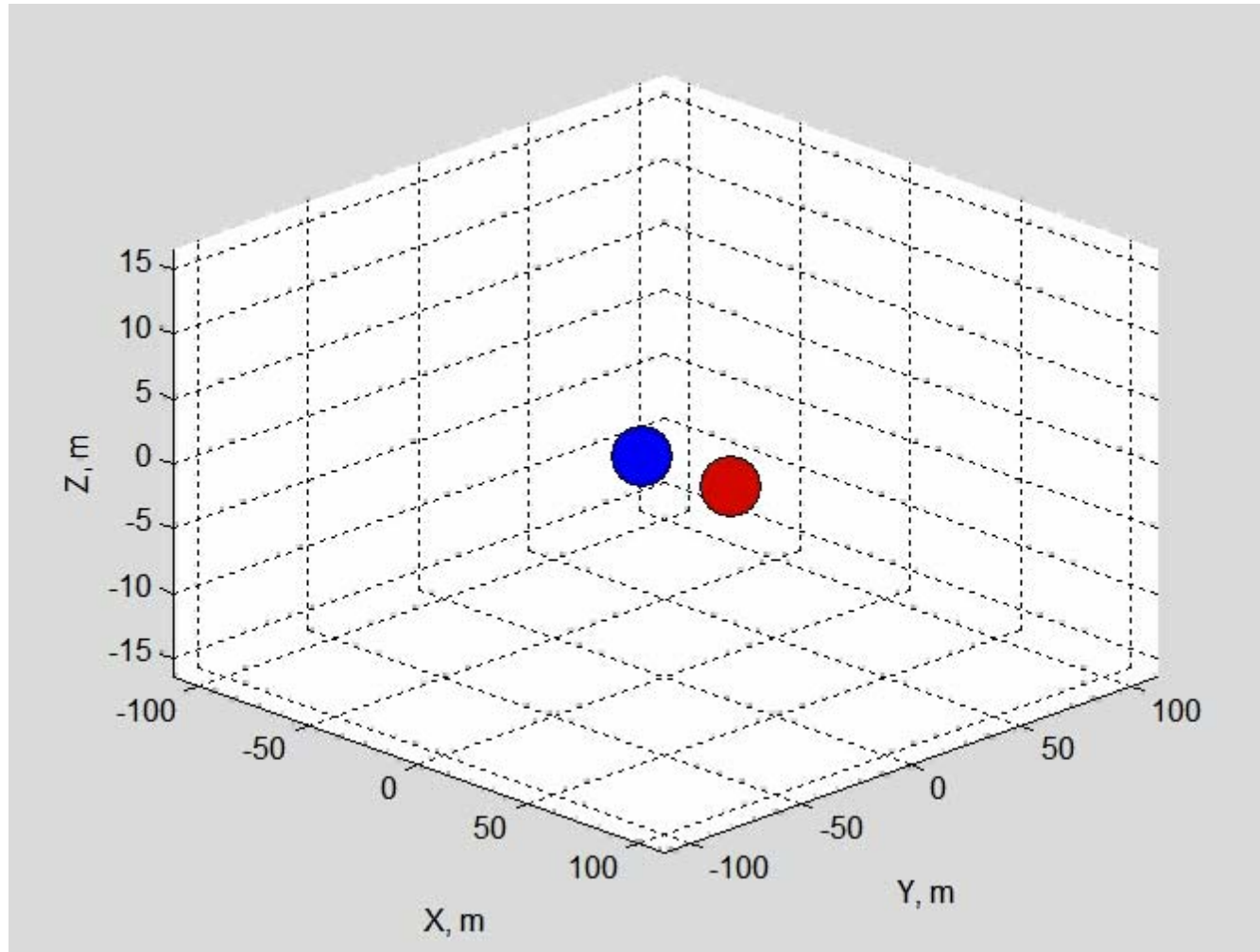
- Example relative orbit
  - S/C 1 Trajectory → like saddle, x-y projection is ellipse



**Full State Periodic Solution: Case B,  $B_z=2$**

# Circumnavigating 2-s/c Orbit

- Another Example
  - Both craft trajectories shown, different mass vehicles



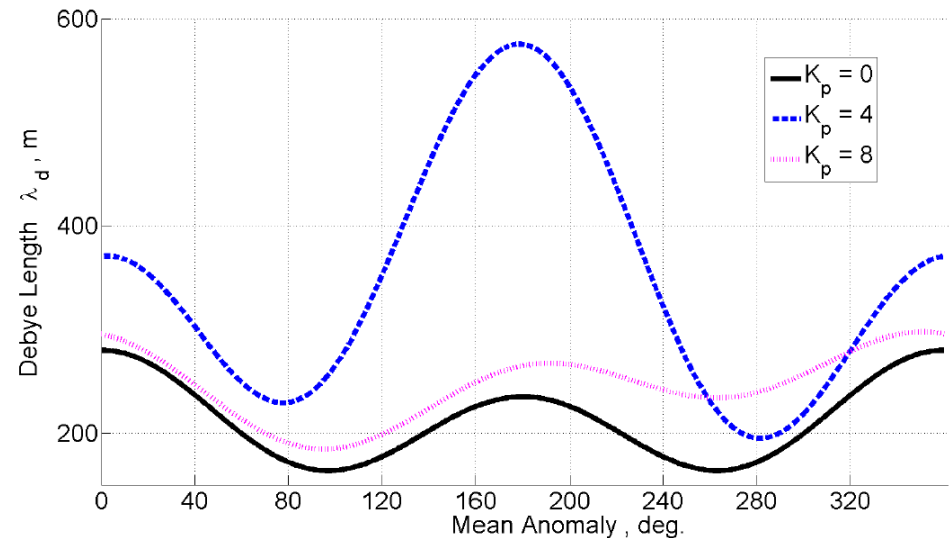
## Perturbed Periodic Coulomb Formation Solutions Propagated in Inertial Frame

- Accuracy of solutions in higher fidelity model
  - Revert back to Newtonian gravity, inertial frame
  - Transformations between ECI and Hill frames
  - Include primary perturbations

## Solar Radiation Pressure using “Cannonball Model”

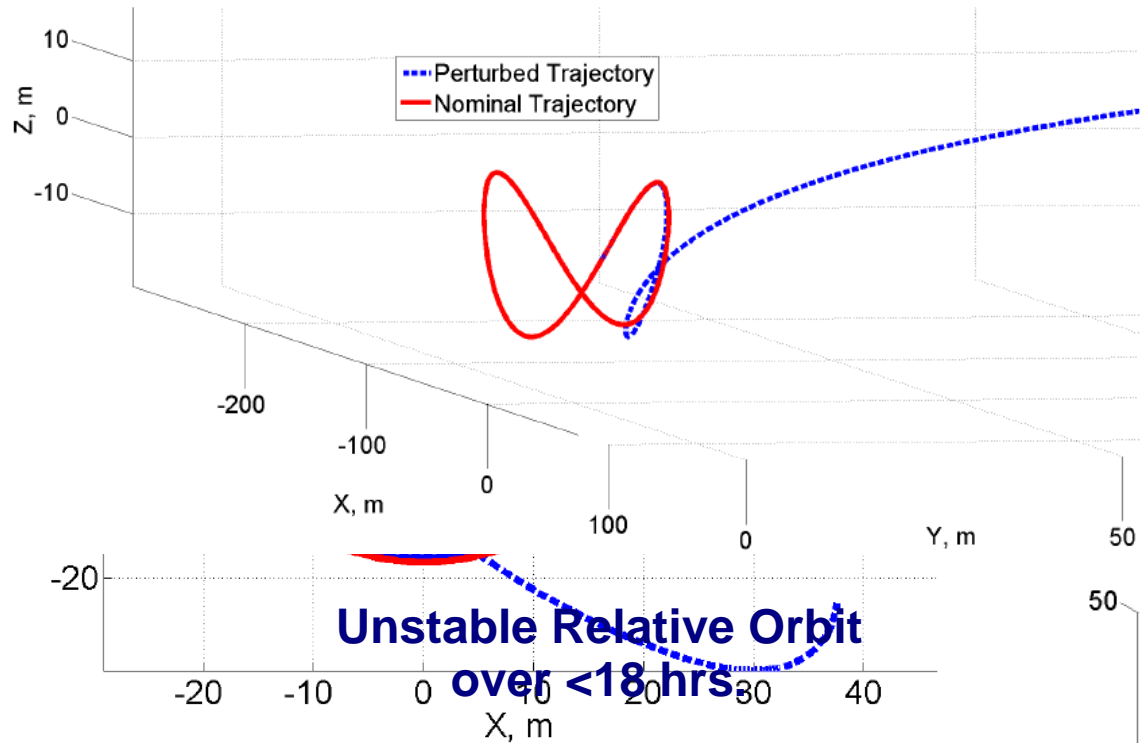
$$f_{\text{srp}} = C_R \frac{\pi R_{sc}^2 \Theta}{m_i c}$$

## Induced Perturbation due to Parametric Uncertainty (Model Error) in Debye Length

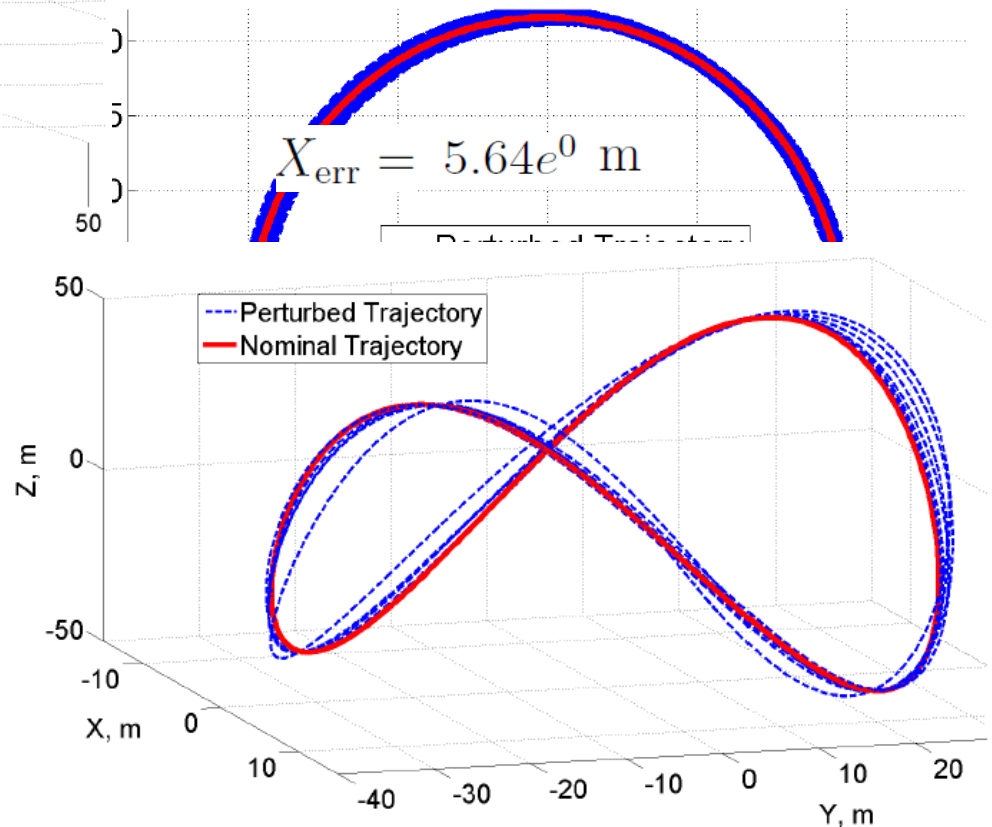




# Perturbed Periodic Solutions



**Unstable Relative Orbit over 10 Revolutions**



- Application of new theory opens multiple pathways
  - The existence of periodic solutions enabled by open-loop electrostatic forcing
  - Stability properties assessed via Floquet theory
  - Some periodic solutions are very nearly stable
- Future Work
  - Other admissible periodic solutions (represent w/ finite Fourier)
  - Feedback control of periodic solutions
  - Test these solutions (and derive new ones) with relaxed assumptions and higher order Coulomb force modeling
  - Precisely measure, control, and estimate s/c potentials (and Coulomb forces) in order to realize these motions in practice?

# QUESTIONS

Please refer further questions to Drew Jones: [drjones604@gmail.com](mailto:drjones604@gmail.com)

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# EXTRA SLIDES

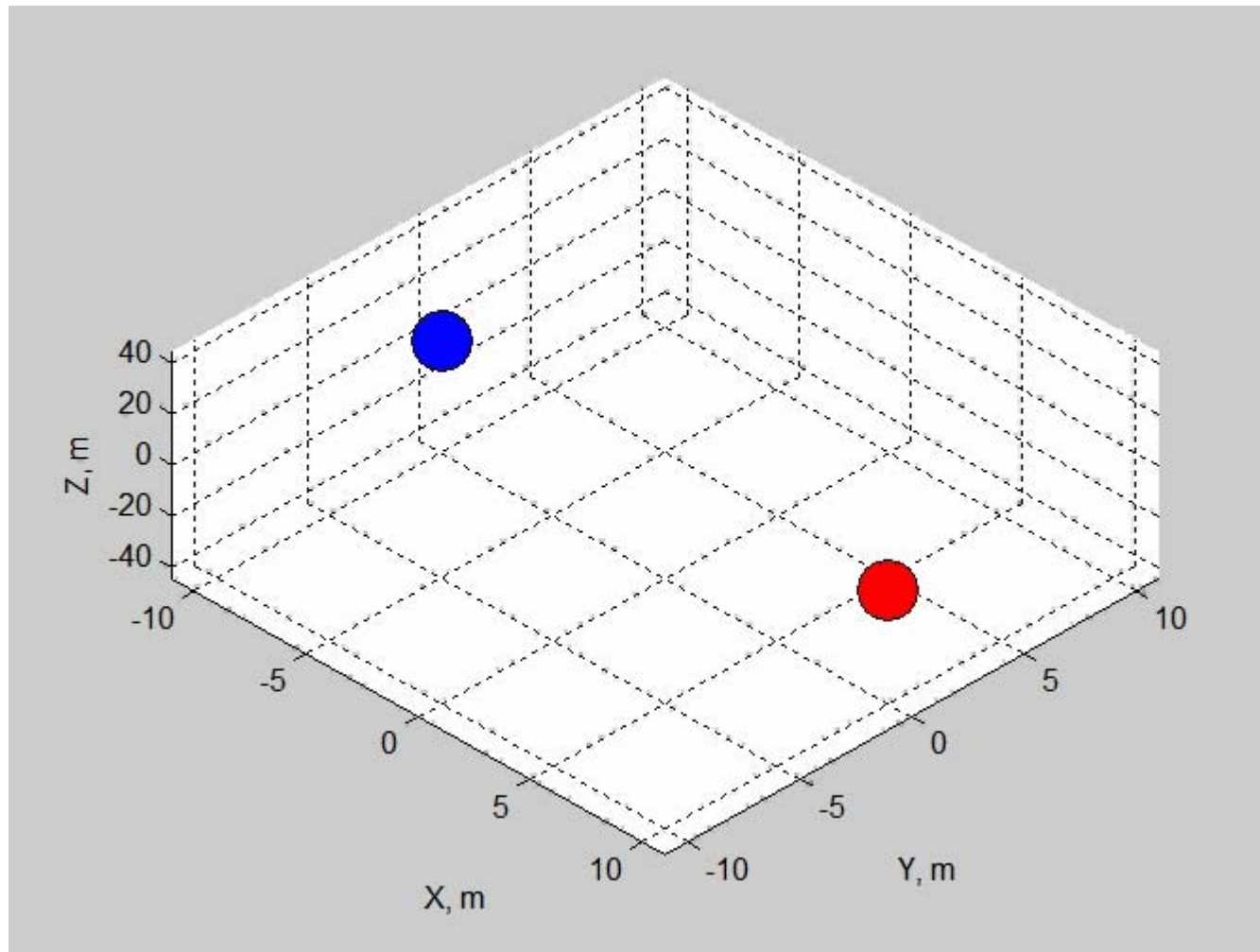
- True force between finite s/c: solution to time-dependent Vlasov-Poisson PDE

$$F \approx k_c \frac{q_1 q_2 (1 + r_{12}/\lambda_d)}{r_{12}^2 \exp[r_{12}/\lambda_d]} \quad \phi_i = k_c \frac{q_i}{R_{sc}}$$

- Assumptions for approximate model accuracy:
  - Craft act equivalently as spheres at a distance
  - Conservative plasma shielding
  - Decoupled capacitances
  - Constant and finite  $\lambda_d$  (at GEO order of 100 m)
  - Formations near GEO ( $\lambda_d$  negligible over  $R_{sc}$ )
  - Lower bound on separation distance:  $r_{12} > 10 R_{sc}$

# Periodic Coulomb Formation

- Case A,  $B_z = 2$ ,  $A_x = 10$  m,  $A_y = 6$  m,  $A_z = 40$  m (17 hrs)



# Periodic Coulomb Formation

- Case B,  $B_z = 4$ ,  $A_x = 20$  m,  $A_y = 125$  m,  $A_z = 5$  m (4 days)

